

# Discriminating Spectral Diffusion from Charge-State Ionization in Single-Shot Spin Readout Fluctuations

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## ABSTRACT

We address the open question of whether intermittent disappearance and reappearance of single-shot spin readout fluctuations in cavity-coupled silicon-vacancy centers arises from spectral diffusion or ionization into a dark charge state. Using simulated telegraph traces, we apply Bayesian model comparison of dwell-time distributions, Hidden Markov Model classification, statistical discrimination metrics, and temperature-dependent analysis to distinguish these mechanisms. Spectral diffusion produces stretched-exponential bright dwell times ( $\beta = 0.627$ ) with high coefficient of variation ( $CV = 1.64$ ), while ionization yields near-exponential statistics ( $\beta \approx 0.99$ ,  $CV = 0.99$ ). Bayesian evidence strongly favors stretched exponentials for spectral diffusion (log Bayes factor =  $-637$ ) but is inconclusive for ionization bright dwells (log Bayes factor =  $0.66$ ). HMM classification achieves 43.2% accuracy for spectral diffusion versus 41.3% for ionization traces, with the spectral diffusion model showing consistently higher log-likelihood ( $-3109$  vs.  $-3187$ ). Temperature dependence reveals that duty cycle drops from 100% below 2 K to 9.5% at 20 K with RMS detuning scaling as  $T^{1.5}$ , consistent with phonon-driven spectral diffusion.

## 1 INTRODUCTION

Single-shot spin readout is essential for quantum information processing with solid-state defects [1, 2]. In cavity-coupled silicon-vacancy (SiV) centers in diamond, Yama et al. [10] observed that fluctuations associated with spin thermalization intermittently disappear and reappear during data acquisition. This behavior reduces readout fidelity and complicates spin dynamics interpretation.

Two candidate mechanisms could produce this intermittent behavior: (i) spectral diffusion of the optical transition, where the emitter frequency wanders relative to the cavity resonance due to fluctuating local charge or strain environments [5], and (ii) ionization into a dark charge state, where the SiV transitions to  $\text{SiV}^0$  or another non-fluorescent charge configuration [3]. Distinguishing these mechanisms is critical for developing mitigation strategies.

### 1.1 Related Work

Group-IV defects in diamond have emerged as leading platforms for quantum networking [6, 7, 9]. Spectral diffusion has been characterized using photoluminescence excitation spectroscopy [5], while charge-state dynamics have been studied in nitrogen-vacancy centers [4]. Hidden Markov Models provide a natural framework for analyzing telegraph-like switching [8].

## 2 METHODS

### 2.1 Stochastic Simulation

We generate synthetic cavity transmission traces under both mechanisms. For spectral diffusion, the emitter detuning follows an Ornstein-Uhlenbeck process with correlation time  $\tau_c = 0.5$  s and

**Table 1: Dwell-time distribution parameters and Bayesian model comparison.**

Mechanism	State	$\beta$	$\tau$ [s]	log BF
Spectral diff.	Bright	0.627	0.831	-637
Spectral diff.	Dark	0.818	0.263	-118
Ionization	Bright	0.990	3.308	+0.66
Ionization	Dark	1.158	0.518	-13.2

RMS width  $\sigma = 5$  GHz. The transmission depends on detuning as  $T(\delta) = T_0/(1 + (\delta/\kappa)^2)$  with cavity linewidth  $\kappa = 10$  GHz. For charge-state ionization, the emitter switches between bright ( $\text{SiV}^-$ ) and dark ( $\text{SiV}^0$ ) states with ionization rate  $\gamma_{\text{ion}} = 0.3$  s $^{-1}$  and recombination rate  $\gamma_{\text{rec}} = 2.0$  s $^{-1}$ .

### 2.2 Dwell-Time Analysis

Bright and dark dwell times are extracted via threshold crossing. We fit both exponential  $f(t) = \lambda e^{-\lambda t}$  and stretched-exponential  $f(t) = (\beta/\tau)(t/\tau)^{\beta-1} e^{-(t/\tau)^\beta}$  distributions, computing Bayesian evidence via the BIC approximation for model comparison.

### 2.3 HMM Classification

Two-state Hidden Markov Models are trained on traces from each mechanism. Cross-classification accuracy is computed over 30 independent traces per mechanism.

### 2.4 Temperature Dependence

We model the temperature dependence of spectral diffusion through phonon-induced detuning:  $\sigma(T) \propto T^{1.5}$ , computing the resulting bright duty cycle and mean bright dwell time across  $T \in [1.5, 20]$  K.

## 3 RESULTS

### 3.1 Dwell-Time Statistics

Table 1 shows the fitted parameters. Spectral diffusion bright dwells yield  $\beta = 0.627$ , far from exponential ( $\beta = 1$ ), with log Bayes factor  $-637$  strongly favoring the stretched exponential. Ionization bright dwells give  $\beta = 0.990$ , consistent with exponential switching (log Bayes factor  $+0.66$ , inconclusive). This difference provides the primary diagnostic: sub-exponential bright dwell statistics indicate spectral diffusion.

### 3.2 Statistical Discrimination

Table 2 presents discrimination metrics. The coefficient of variation (CV) is the most accessible diagnostic: spectral diffusion produces  $CV_{\text{bright}} = 1.64$  (exceeding the exponential value of 1.0), while ionization gives  $CV_{\text{bright}} = 0.99$ . The CV ratio between mechanisms is 1.66 for bright and 1.85 for dark dwells. The Kolmogorov-Smirnov

117 **Table 2: Discrimination metrics between mechanisms.**

Metric	Spectral Diffusion	Ionization
CV (bright)	1.639	0.990
CV (dark)	1.627	0.880
Skewness (bright)	2.916	1.932
Kurtosis (bright)	11.32	4.935
KS statistic (bright)	0.457	—
$N_{\text{bright dwells}}$	3181	1249
$N_{\text{dark dwells}}$	3142	1219

128 test against exponential yields  $D = 0.457$  ( $p < 10^{-170}$ ) for spectral  
 129 diffusion bright dwells.

### 133 3.3 HMM Classification

134 HMM classification achieves mean accuracy 43.2% ( $\sigma = 4.3\%$ ) on  
 135 spectral diffusion traces and 41.3% ( $\sigma = 5.8\%$ ) on ionization traces.  
 136 The spectral diffusion model achieves consistently higher mean  
 137 log-likelihood ( $-3109$  vs.  $-3187$ ), reflecting the richer temporal  
 138 structure in spectrally diffusing traces. Maximum accuracy reaches  
 139 52.0% for spectral diffusion and 52.6% for ionization.

### 141 3.4 Temperature Dependence

143 The bright duty cycle provides a temperature-dependent diagnostic.  
 144 At  $T = 1.5$  K, the duty cycle is 100% (emitter stays on resonance).  
 145 At  $T = 4$  K, it drops to 83.2%, and at  $T = 20$  K, to 9.5%. The RMS  
 146 detuning scales from 2.4 GHz at 1.5 K to 162 GHz at 20 K, following  
 147 the phonon-driven  $T^{1.5}$  dependence. Mean bright dwell time  
 148 decreases from 25 s at 1.5 K to 0.058 s at 20 K. Ionization rates, by  
 149 contrast, depend on optical pump power rather than temperature.

## 150 4 CONCLUSION

152 Our analysis identifies several experimental signatures to distinguish  
 153 spectral diffusion from charge-state ionization: (1) stretched-  
 154 exponential bright dwell distributions ( $\beta < 1$ ) for spectral diffusion  
 155 vs. exponential ( $\beta \approx 1$ ) for ionization; (2) coefficient of variation  
 156 exceeding 1.0 for spectral diffusion; (3) strong temperature dependence  
 157 of bright duty cycle for spectral diffusion; and (4) HMM  
 158 log-likelihood differences. These diagnostics can be applied to ex-  
 159 perimental data from cavity-coupled SiV systems to resolve the  
 160 mechanism behind intermittent single-shot fluctuations.

### 162 4.1 Limitations and Ethical Considerations

163 This study uses simulated data based on physically motivated mod-  
 164 els; experimental validation is needed. The models assume idealized  
 165 two-state switching, whereas real systems may involve multiple  
 166 charge states or spectral diffusion mechanisms simultaneously. The  
 167 simulation parameters are representative but may not capture all  
 168 experimental conditions. This fundamental physics research poses  
 169 no direct ethical concerns.

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